

Study of the dynamics of hydrogen bonds in water and consequences for the unusual behaviour of supercooled water

José Teixeira and Stéphane Longeville

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Liquid Water

supercooling

apparent divergence
of thermal properties

Lowest temperature of
available data ~ 240 K

exp. $T_H = 231$ K
(1958)

Extrapolation/
Speculation
 $T_S = 228$ K
(1976)

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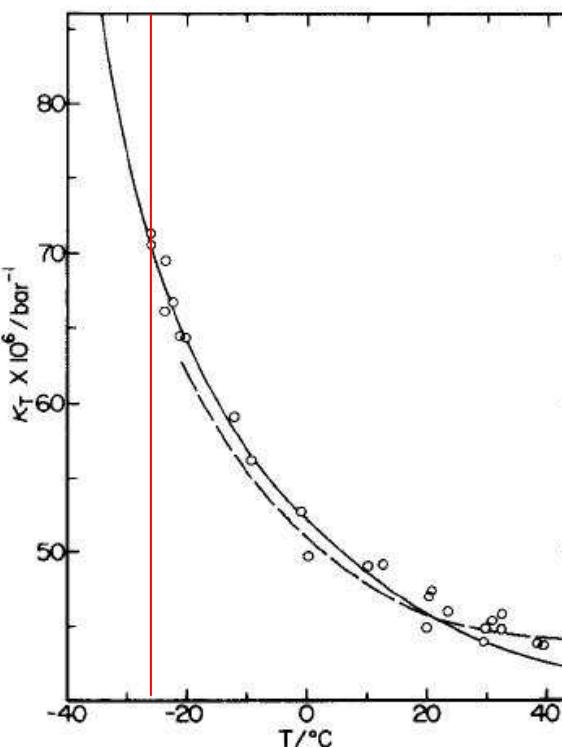
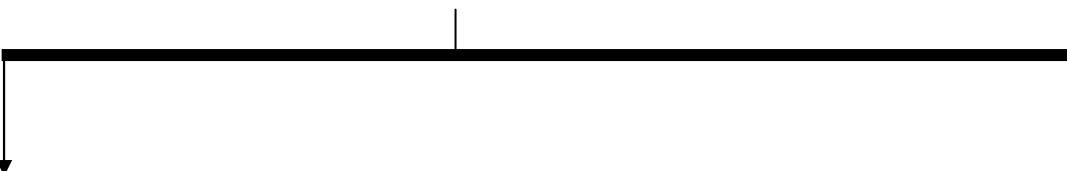


FIG. 2. Compressibility of water in normal and supercooled states. Dashed line is polynomial extrapolation by Kell, (Ref. 13) of data from above 0°C. Solid line is plot of Eq. (2) for parameters given in text.

R. Speedy and C.A. Angell
J. Chem. Phys. **65**, 851 (1976)

Liquid Water



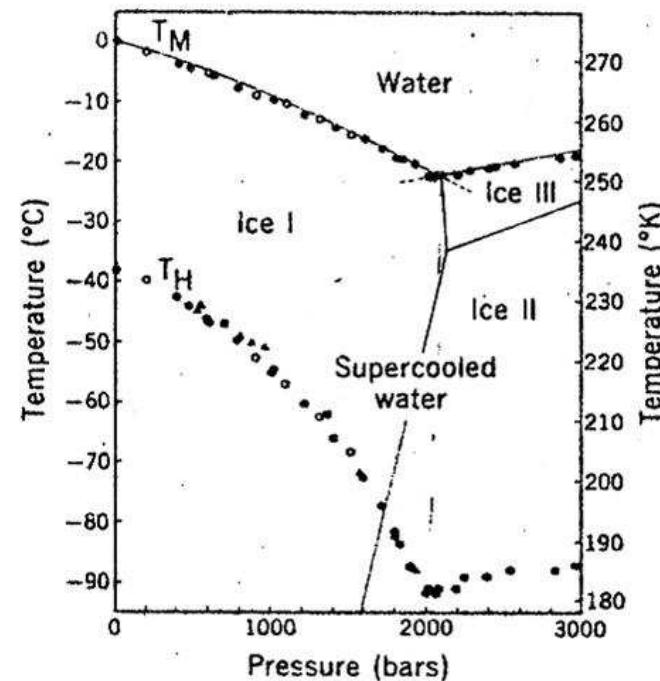
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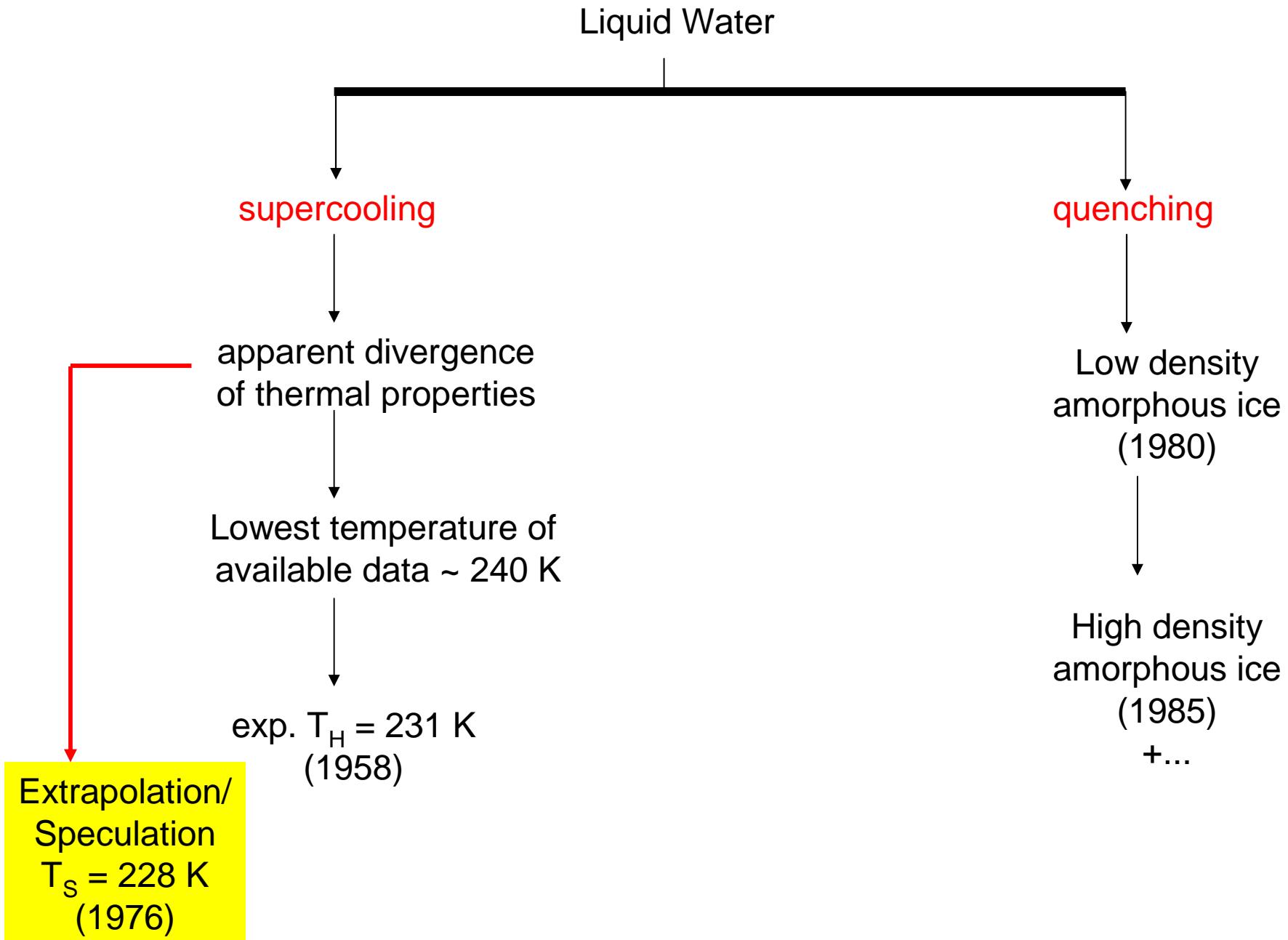
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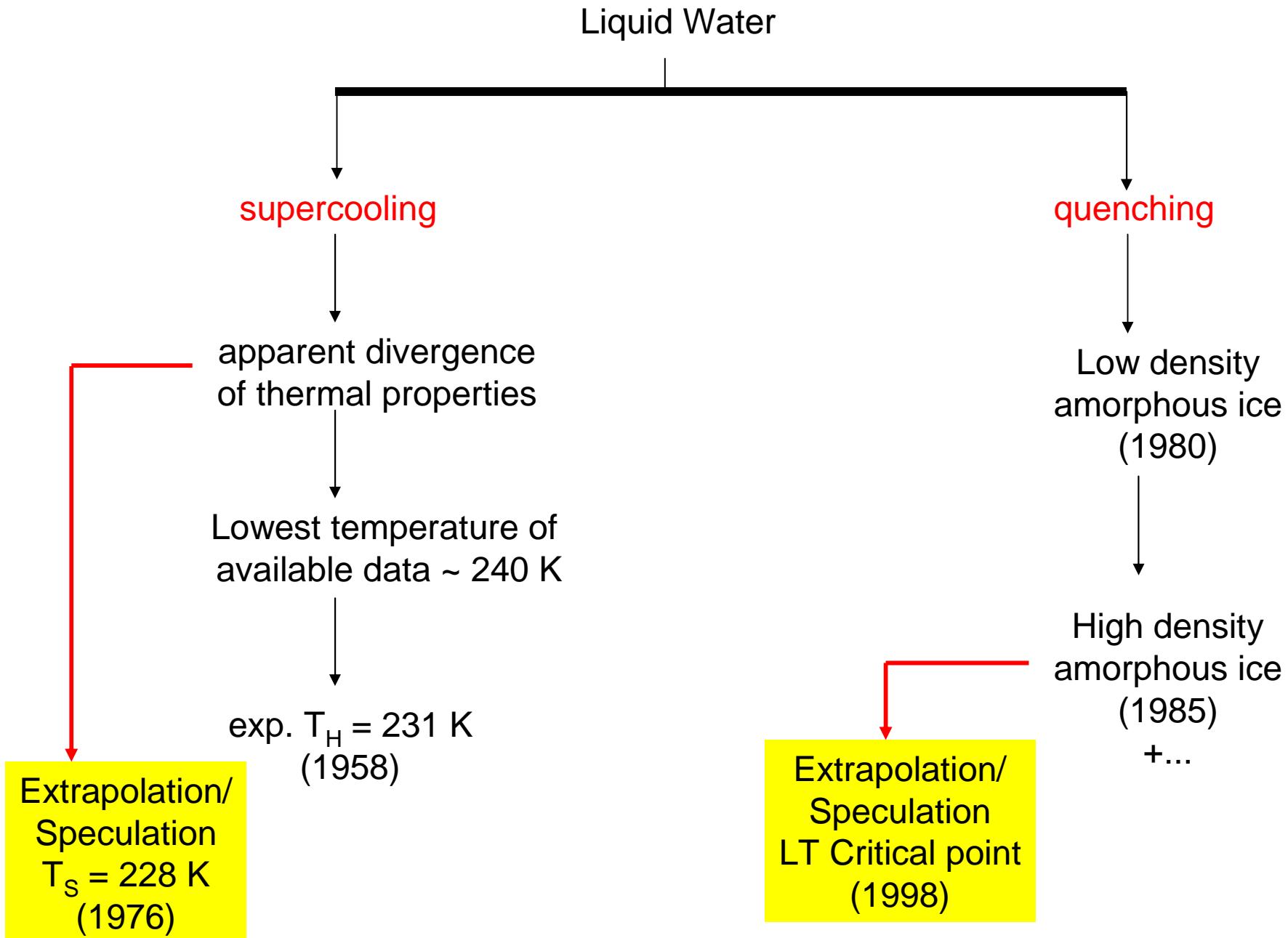
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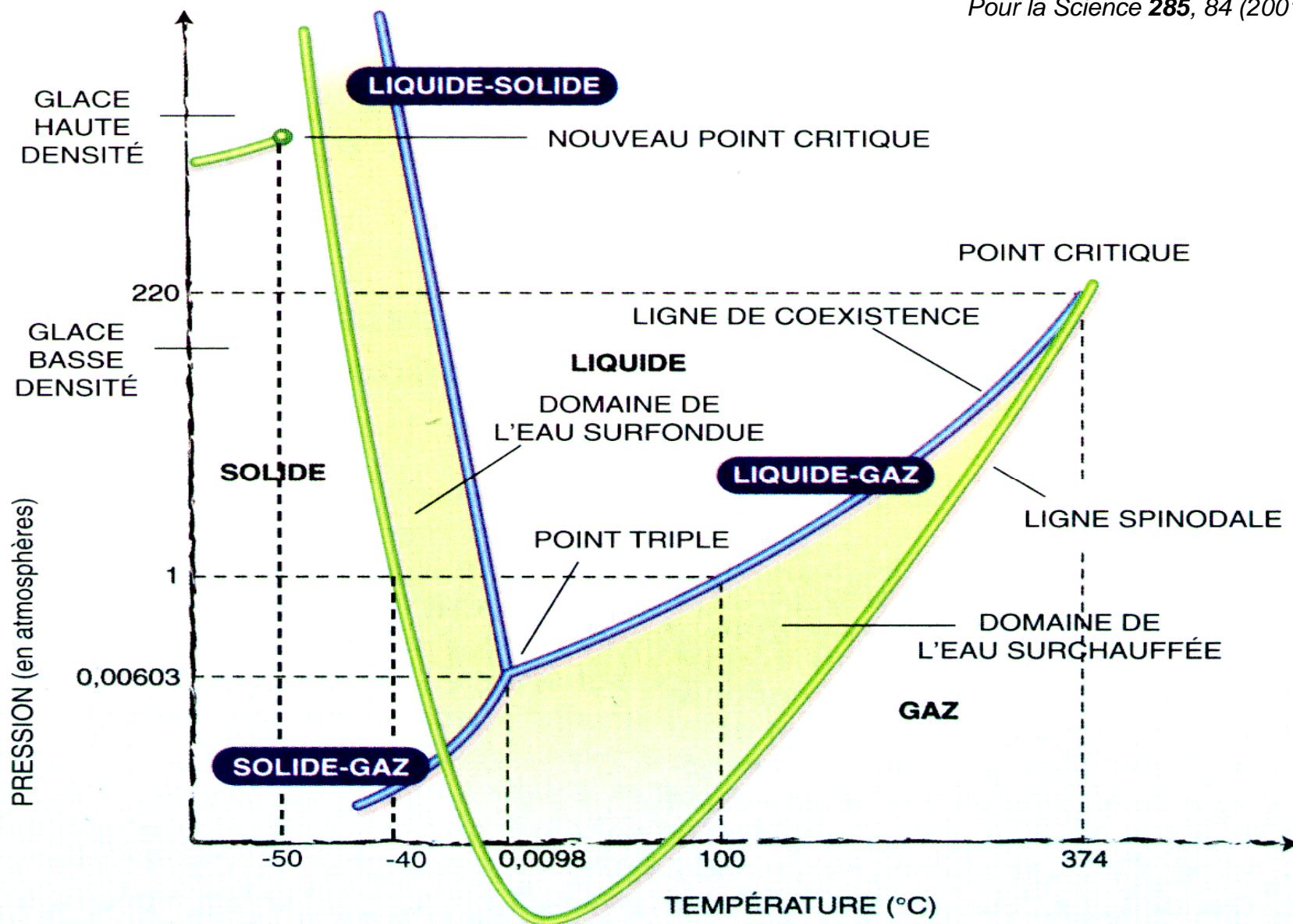
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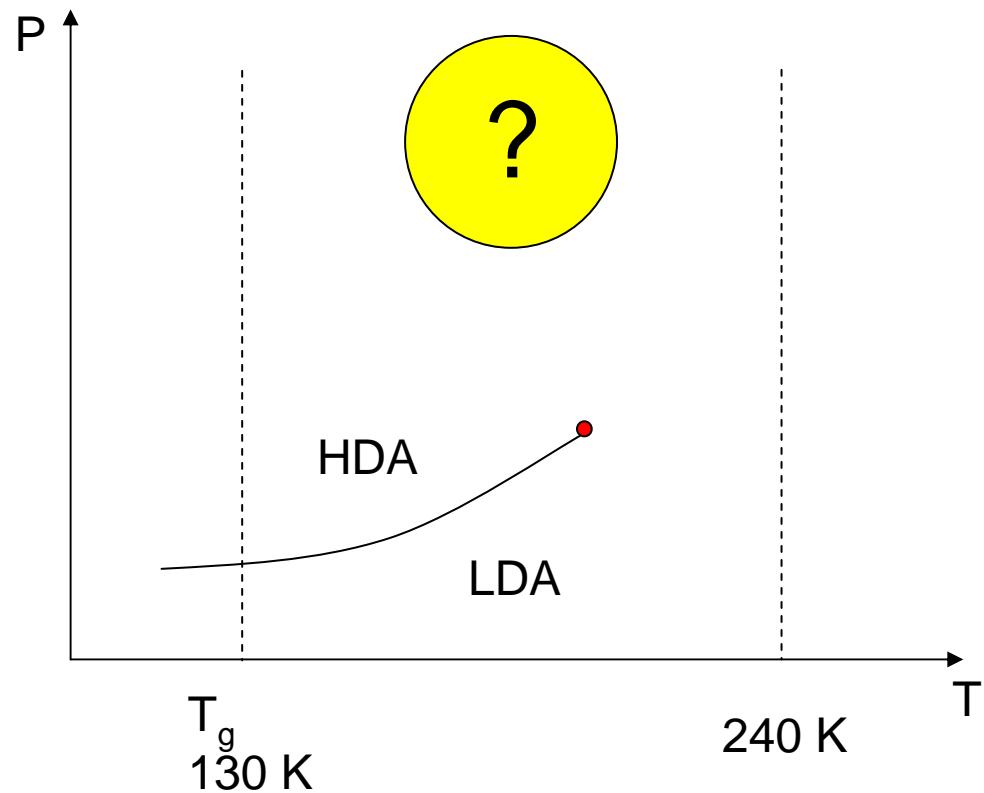


H. Kanno, R.J. Speedy and C.A. Angell
Science 189, 880 (1975)

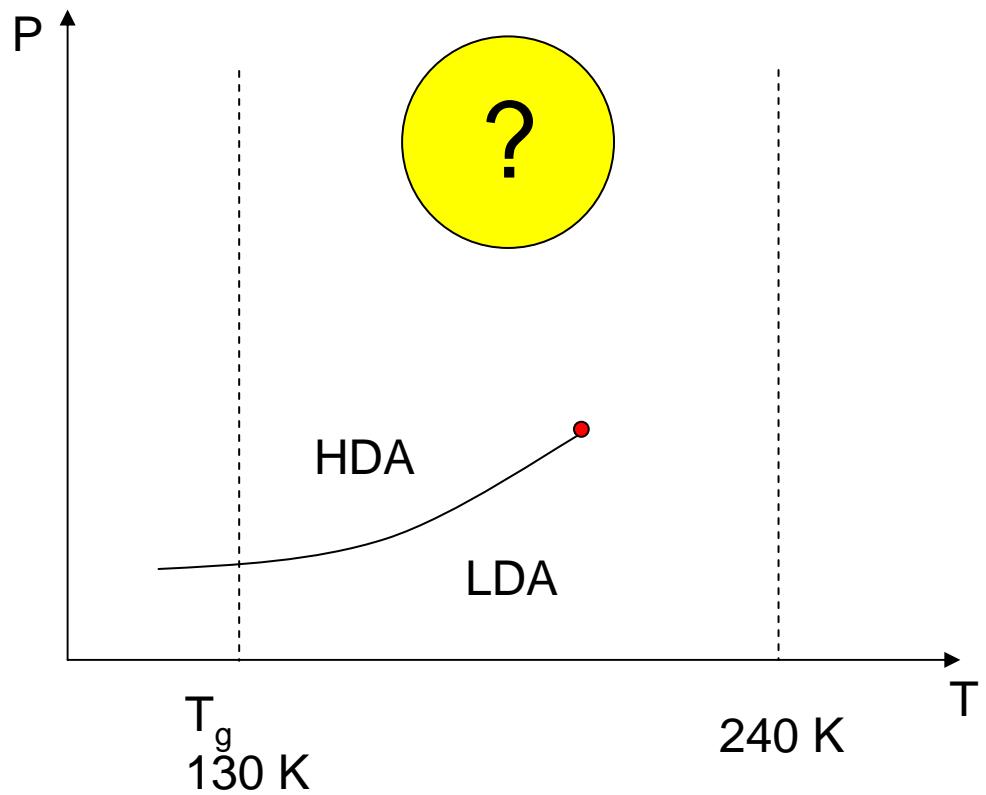








O. Mishima and H.E. Stanley
Nature 396, 329 (1998)



Computer simulations
of
molecular dynamics
using
« effective potentials »

Acronym (date)	Status	Type	Sites	$\mu_g(D)$	$\mu_t(D)$
BF (1933)	empirical	R	4	2.0	2.0
R (1951)	empirical	R	5	1.84	1.84
BNS (1971)	empirical	R	5	2.17	2.17
ST2 (1973,1993)	empirical	R	5	2.35	2.35
CF (1975,1978,1995)	empirical	F	3	1.86	1.98
MCY (1976)	ab initio	R	4	2.19	2.19
DCF (1978,1980,1993)	empirical	F,D,P	3	1.855	-
PE (1979)	empirical	P	1	1.855	2.50
SPC (1981)	empirical	R	3	2.27	2.27
TIP3P (1981,1983)	empirical	R	3	2.35	2.35
RWK (1982)	empirical	F	4	1.85	1.89
TIP4P (1983)	empirical	R	4	2.18	2.18
BJH (1983)	empirical	F	3	1.87	1.99
SPC/F (1985)	empirical	F	3	2.27	2.42
MCYL (1986)	ab initio	F	4	2.19	2.26
SPC/E (1987)	empirical	R	3	2.35	2.35
WK (1989)	empirical	R	4	2.60	2.60
SPCP (1989)	empirical	P	3	1.85	2.90
CKL (1990)	empirical	F,P	4	1.88	2.20
MCHO (1990)	ab initio	P	6	2.12	≈ 3.0
NCC (1990)	ab initio	P	6	1.85	2.80
NEMO (1990,1995)	ab initio	P	5	2.04	2.89
PTIP4P (1991)	empirical	P	4	1.85	2.80
SPC/FP (1991)	empirical	F,P	3	1.85	2.44
PSRWK (1991)	empirical	P	4	1.88	2.63
KJ (1992)	empirical	P	4	1.85	-
NCCvib (1992)	ab initio	F,P	6	1.85	3.11
ASP-W (1992,1998)	ab initio	P	3	1.85	2.90
RPOL (1992)	empirical	P	3	2.02	2.62
CPMD (1993,1999)	DFT+CP	F,D,P	nucl.+el.	1.87	2.95
PPC (1994)	ab initio	P	4	2.14	2.51
SPC/FQ (1994)	empirical	P	3	1.85	2.83
TIP4P/FQ (1994)	empirical	P	4	1.85	2.62
KKY (1994)	empirical	F,D	3	2.38	2.21
SQPM (1995)	valence bond	P	4	1.85	2.62
SCPDP (1996)	empirical	P	4	1.85	2.87
TAB/10D (1998)	SCF+MD	P	5	1.85	2.65
NSPCE (1998)	empirical	R	3	2.18	2.18
NCF (1998)	empirical	F	3	1.85	1.90
MCDHO (2000)	ab initio	F,P	4	1.85	3.01
TIP5P (2000)	empirical	R	5	2.29	2.29
SPC/HW (2001)	empirical	R	3	2.41	2.41
DEC (2001)	empirical	R	3	1.85	1.85
SWFLEX (2001)	empirical	P	4	1.85	2.59
POLARFLEX (2001,2003)	valence bond	F,P	3	1.85	2.55,2.59
POL5 (2001)	ab initio	P	5	1.85	2.71
TTM2/R,F (2002,2006))	ab initio	R(F),P	4	1.85	2.67
SPC/S (2002)	empirical	R	3	2.30	2.30
COS/B2(2002)	empirical	P	3	2.07	2.62
AMOEBA(2003)	empirical	F,P	3	1.77	2.78
SWM4-DP (2003)	empirical	P	4	1.85	2.46
TIP5P/E (2004)	empirical	R	5	2.29	2.29
ABEEM-7P (2004)	empirical	F,P	7	1.85	2.80
COS/G2,G3 (2004)	empirical	P	4	1.85	2.59
TIP3P-PME (2004)	empirical	R	3	2.34	2.34
VRT(ASP-W)III (2005)	ab initio+VRT	F,P	3	1.85	-
GCPM (2005)	empirical	P	4	1.85	2.72
TIP4P/Ew,2005,Ice (2005)	empirical	R	4	2.321, 2.305, 2.426	
QMFF (2006)	ab initio	F,P	6	1.87	-

Bertrand Guillot
2006

Good thermal properties → « Excess » of structure

Good structure → Less « anomalies »

Parameters optimized for room temperature liquid water
Quantum effects not always taken into account

...

Main problem:

Simulations

Point charges

Spherical (electrostatic) potentials

Real water

Spherical molecule

Anisotropic potential



Hydrogen bonds

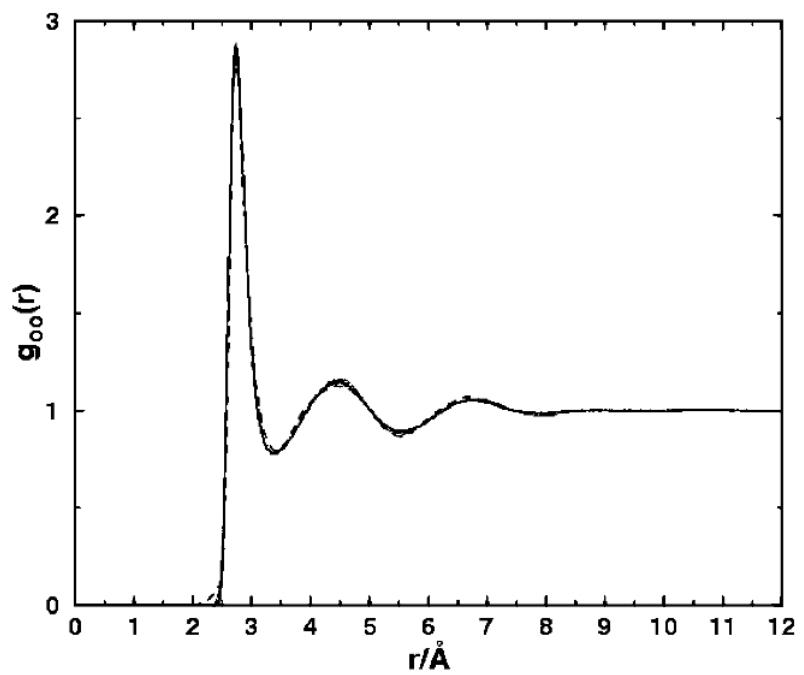
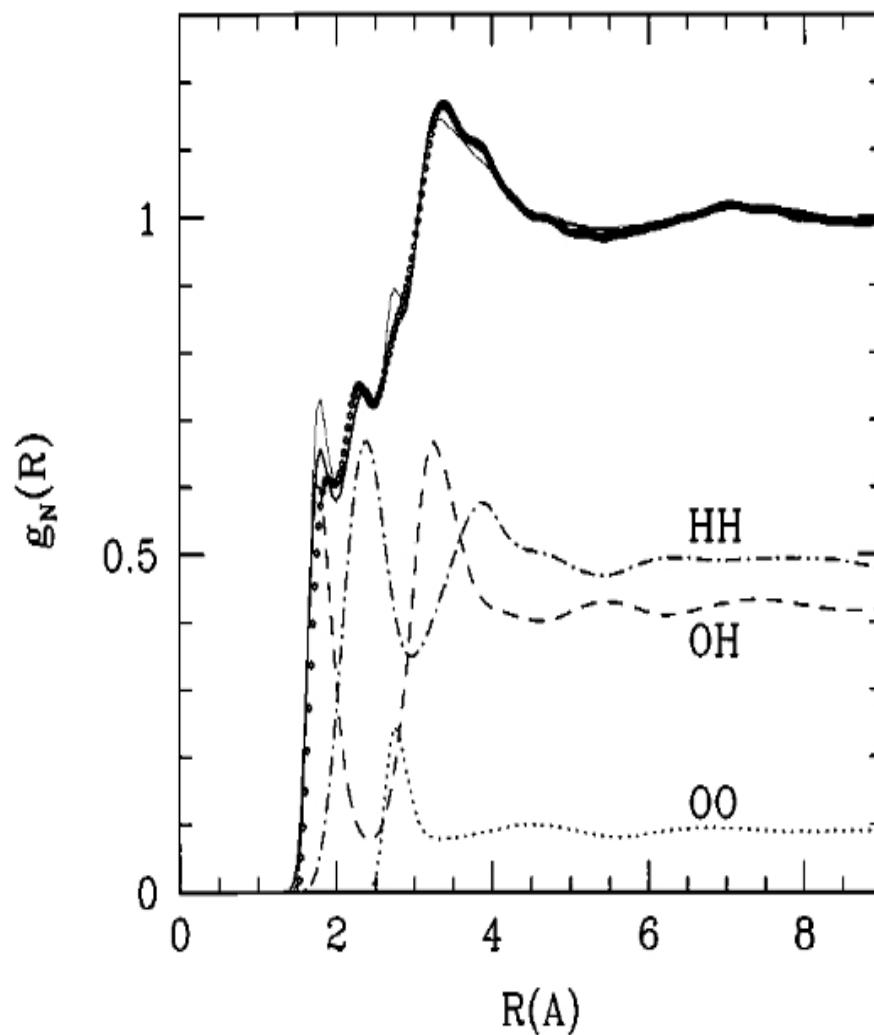
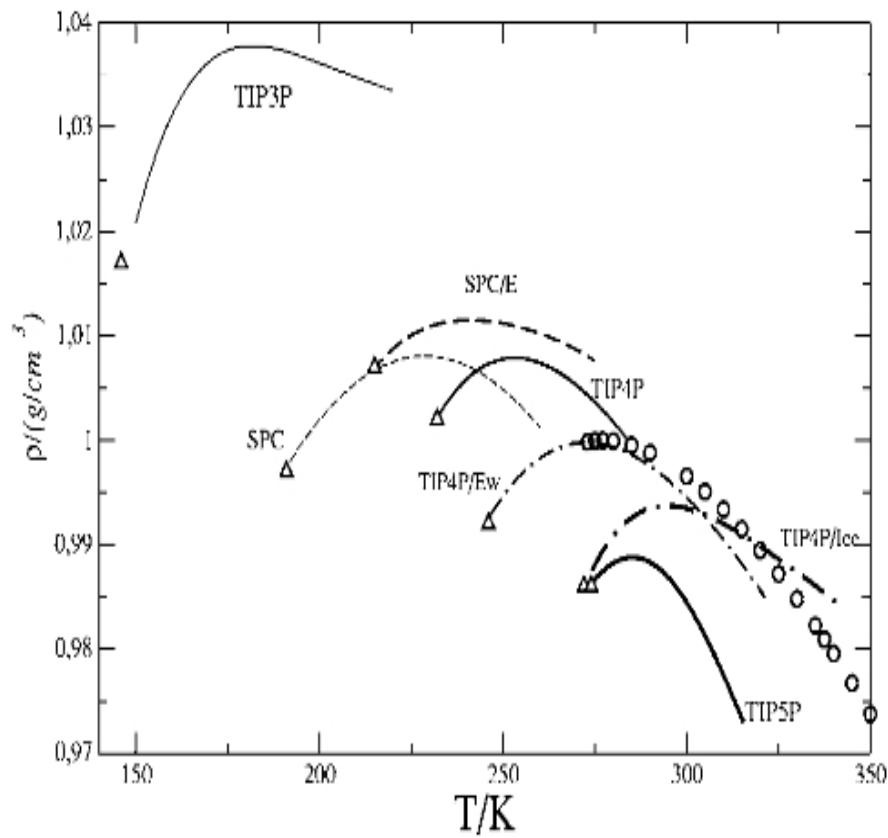


Fig. 1 Comparison of $g_{OO}(r)$ derived from r -space fitting using ALS experimental data² (dashed line) and neutron experimental data and empirical potential structural refinement (black line).¹⁴



RX: Hura et al. (2000), JCP 113, 9140
 Neutrons: Soper (2000), Chem. Phys. 258, 121

Simulation: Guillot and Guissani (2001), JCP 114, 6720
 Exp.: Bellissent-Funel et al. (1997), JCP 107, 2942



Vega et al. (2005), JCP **123**, 144504

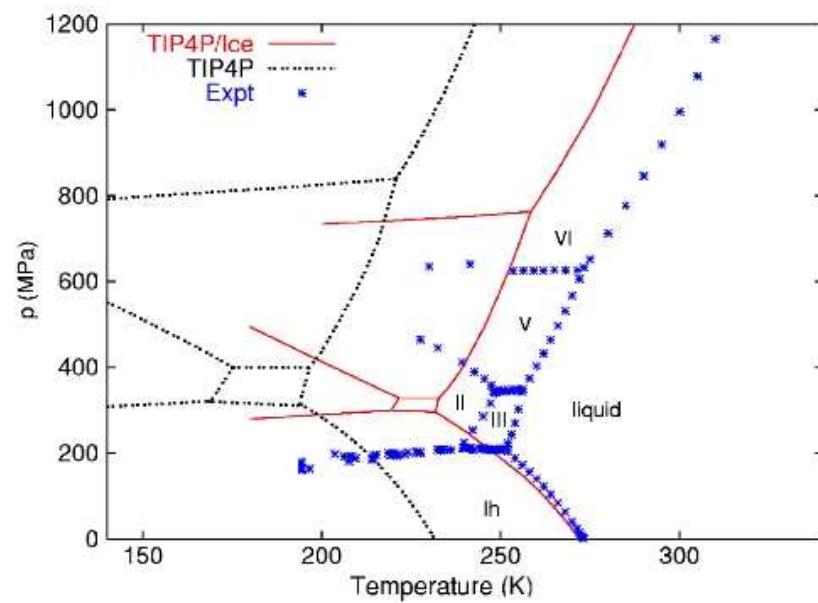
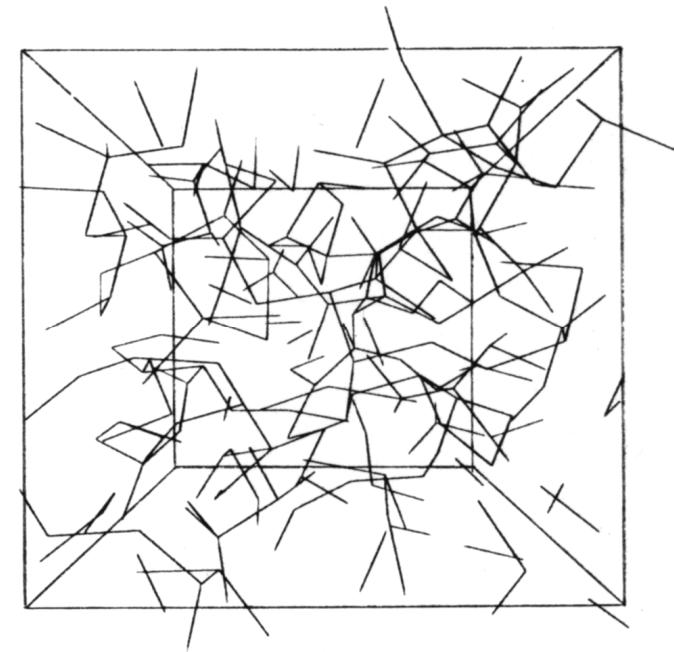


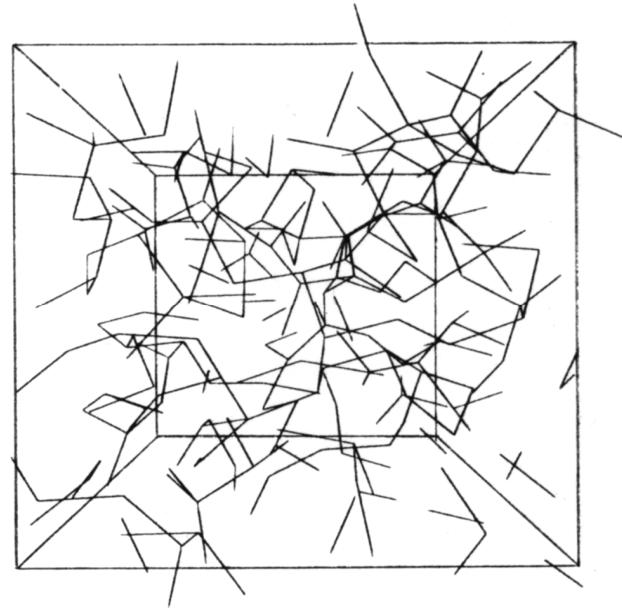
FIG. 1. The phase diagram of TIP4P/Ice (full lines) and TIP4P (dashed) compared to the experiment (stars). The labels mark the domain of stability of the ice phases in the experimental phase diagram.

Abascal et al. (2005), JCP **122**, 234511

- Large number of hydrogen bonds
- Local tetrahedral symmetry
- Short life time of the bonds

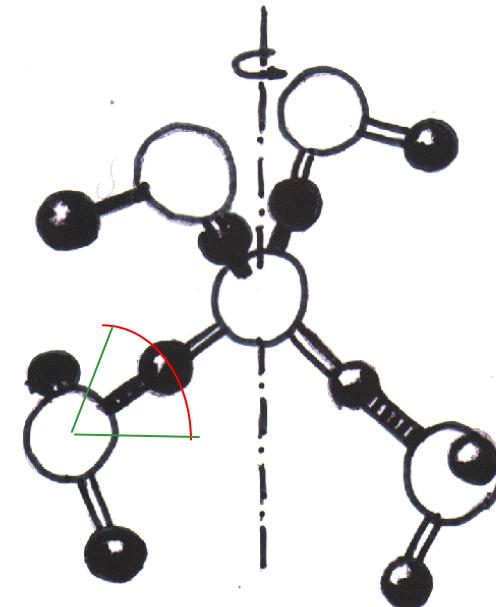
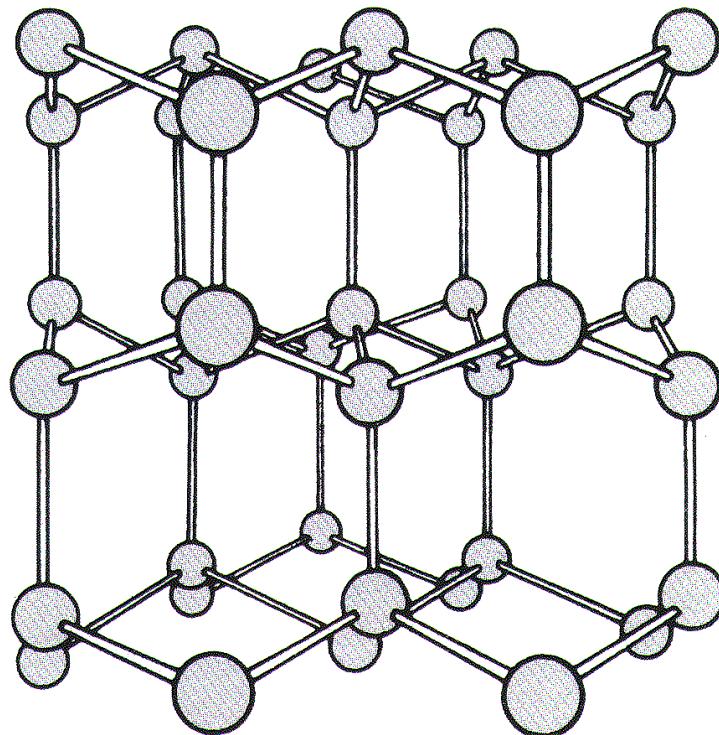


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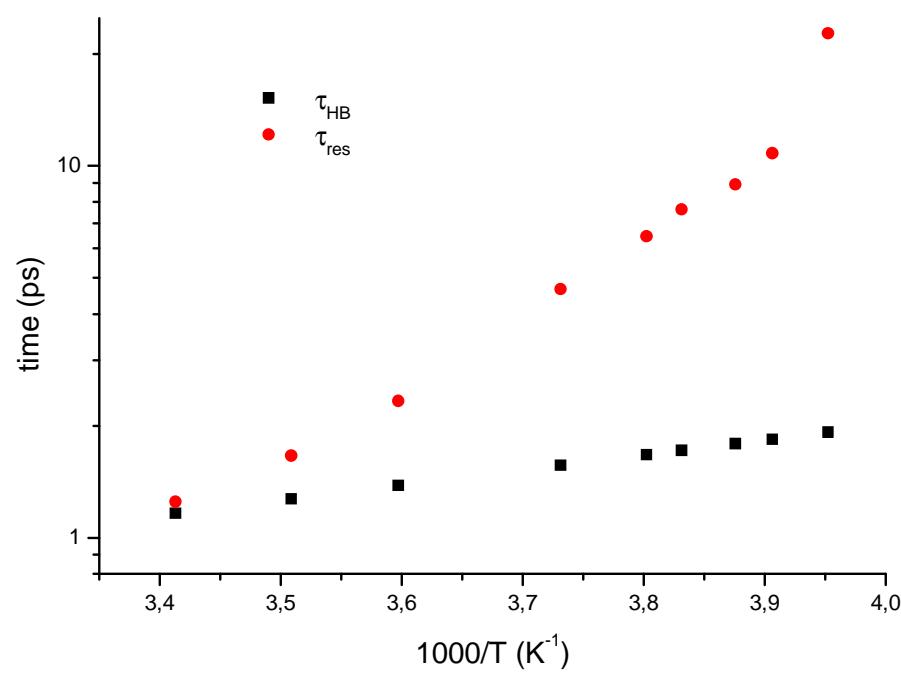
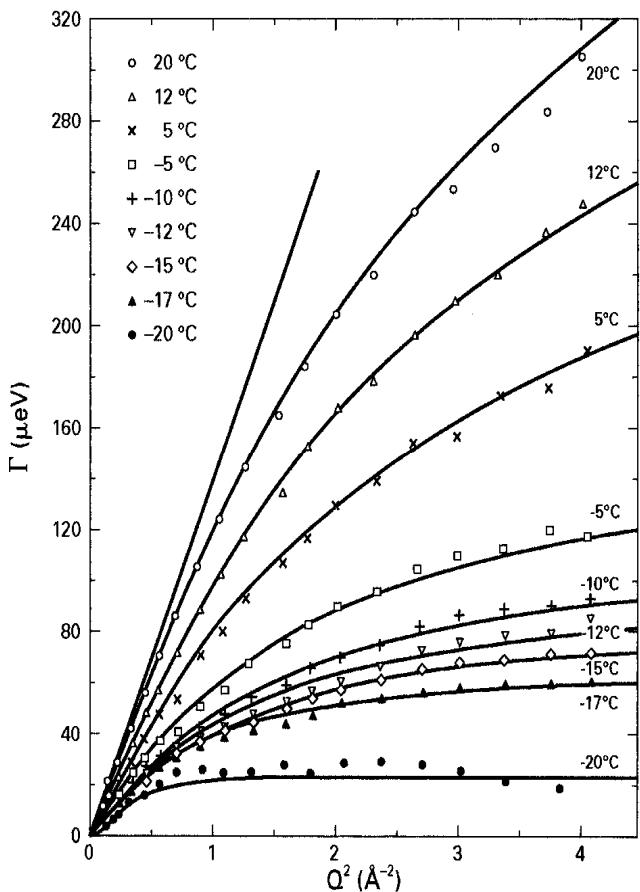


How to tackle HB dynamics?

Librations and HB dynamics



HB and residence times (QENS)



*J. Teixeira, M.-C. Bellissent-Funel,
S.-H. Chen and A.J. Dianoux
Phys. Rev. A31, 1913 (1985)*

Bulk and confined water

$$\tau_{HB} / ps = 0.0485 \exp(931/T)$$

$$E_{HB} = 1.85 \text{ kcal/mol}$$

$$p_B = n_{HB} / 4 = 1.8 - 0.004T$$

Bulk

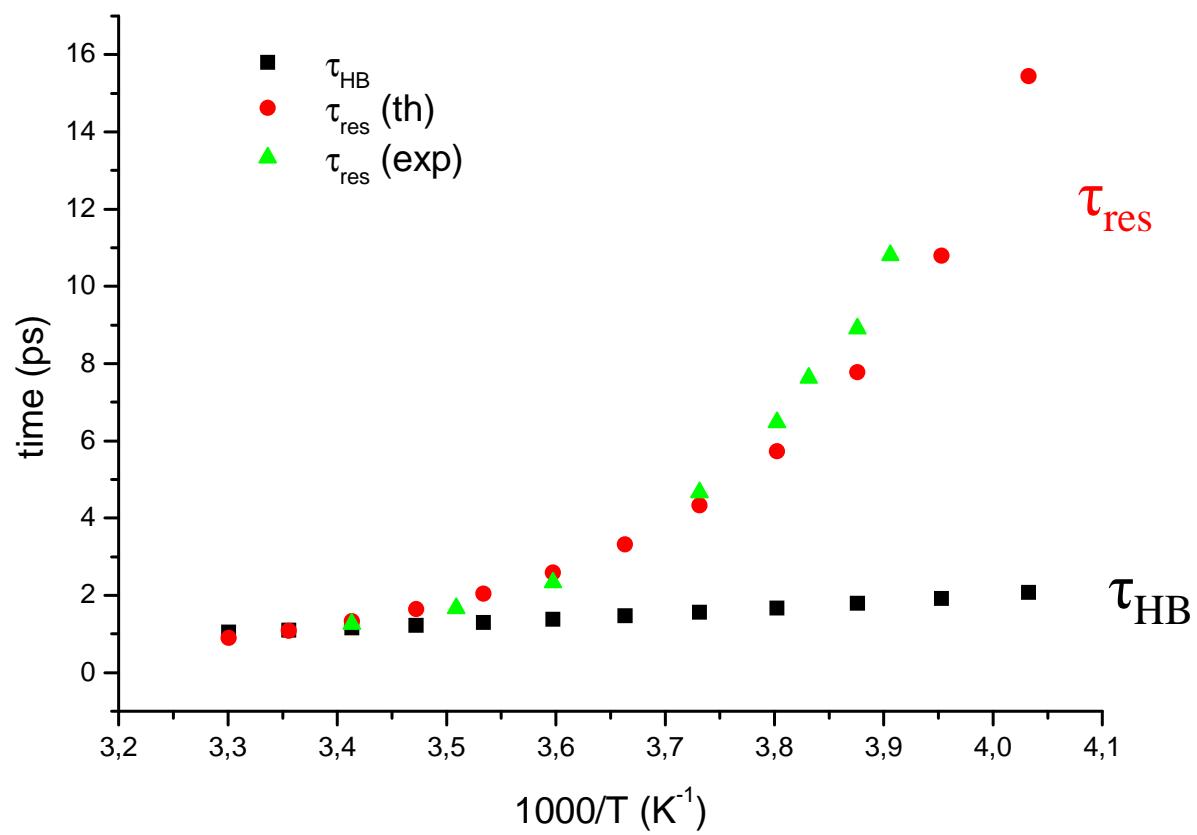
$$F_I^4 = 6p_B^2 - 8p_B^3 + 3p_B^4$$

$$\tau_{res}^{bulk} = const \left(\frac{\tau_{HB}}{-\ln(F_I^4)} \right)$$

Confined

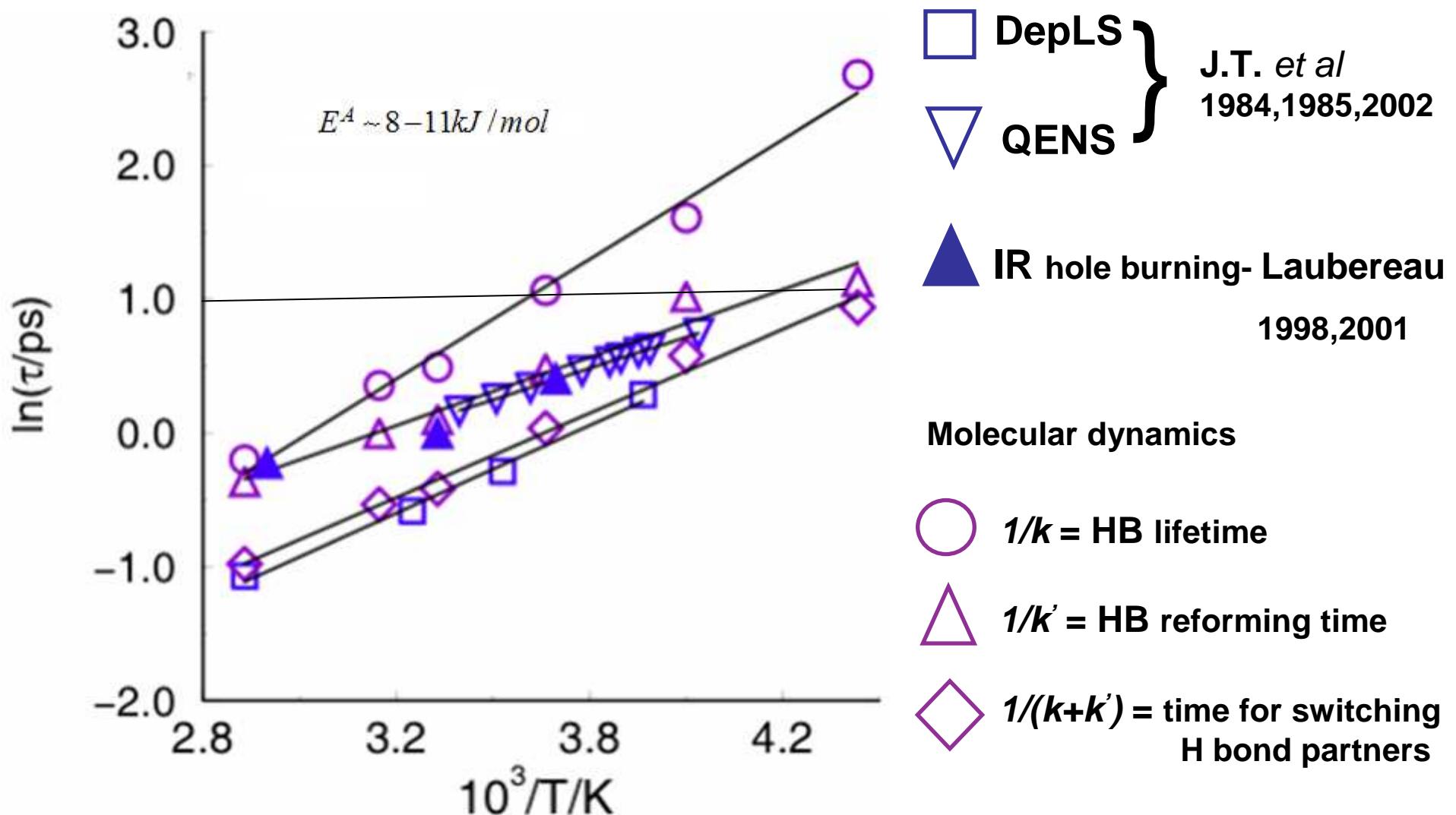
$$F_I^3 = 3p_B - 3p_B^2 + p_B^3$$

$$\tau_{res}^{conf} = const \left(\frac{\tau_{HB}}{-\ln(F_I^3)} \right)$$



Microscopic measurements of the dynamics of H bonds show there are NO anomalies at their level

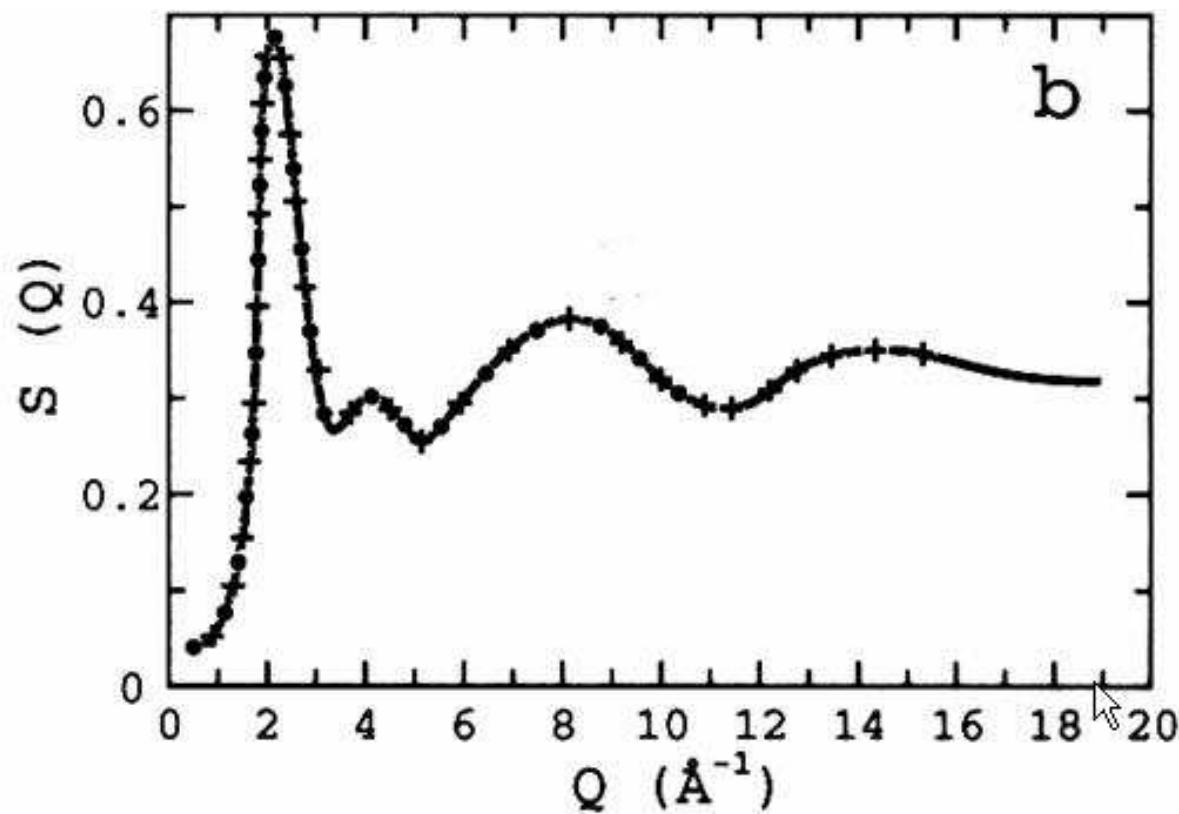
J. Teixeira, A. Luzar and S. Longeville, J. Phys.: Cond. Matter **18**, S2353 (2006)



Microscopic measurements of the dynamics of H bonds show there are
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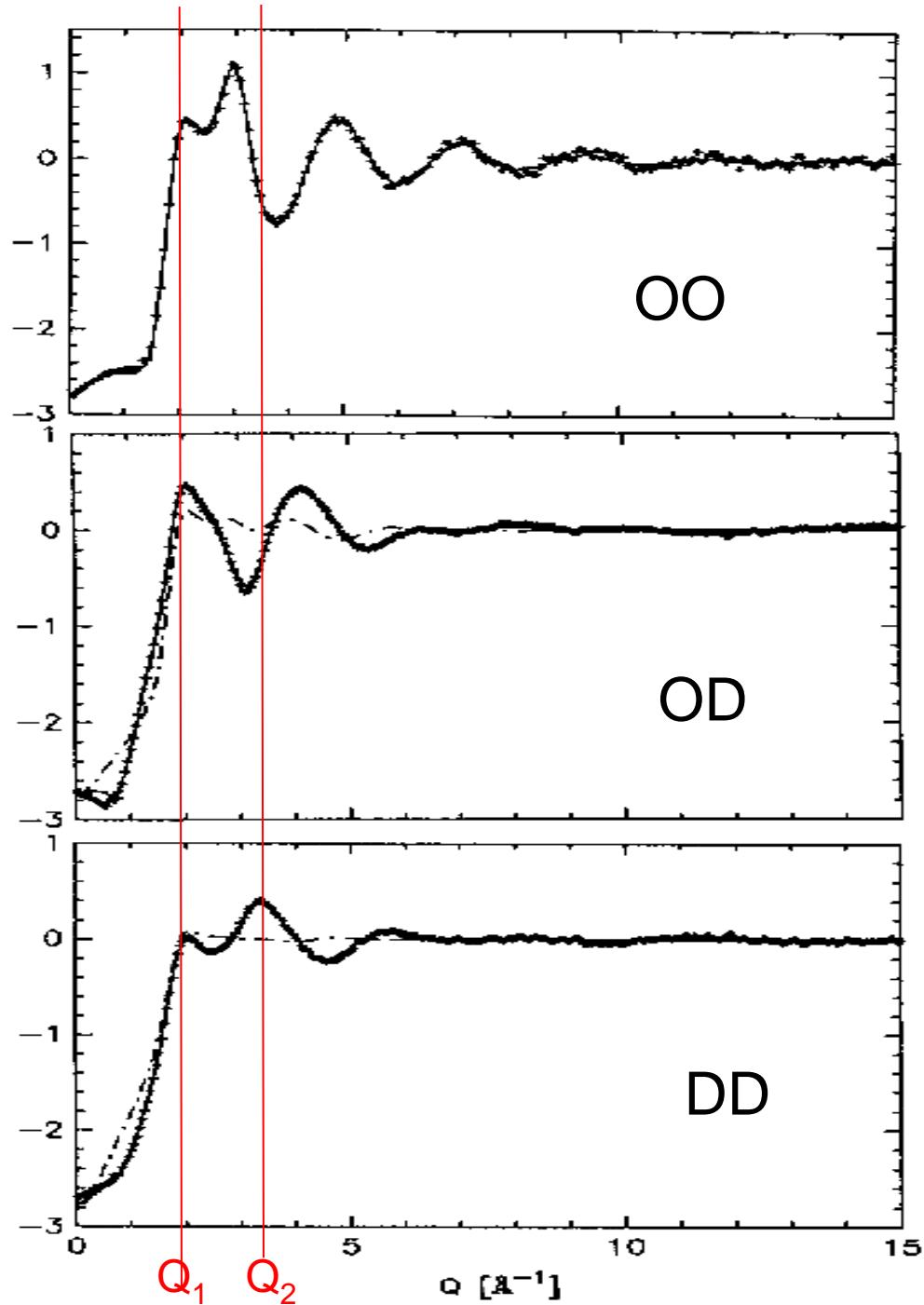
Structure factor of liquid water (neutrons)



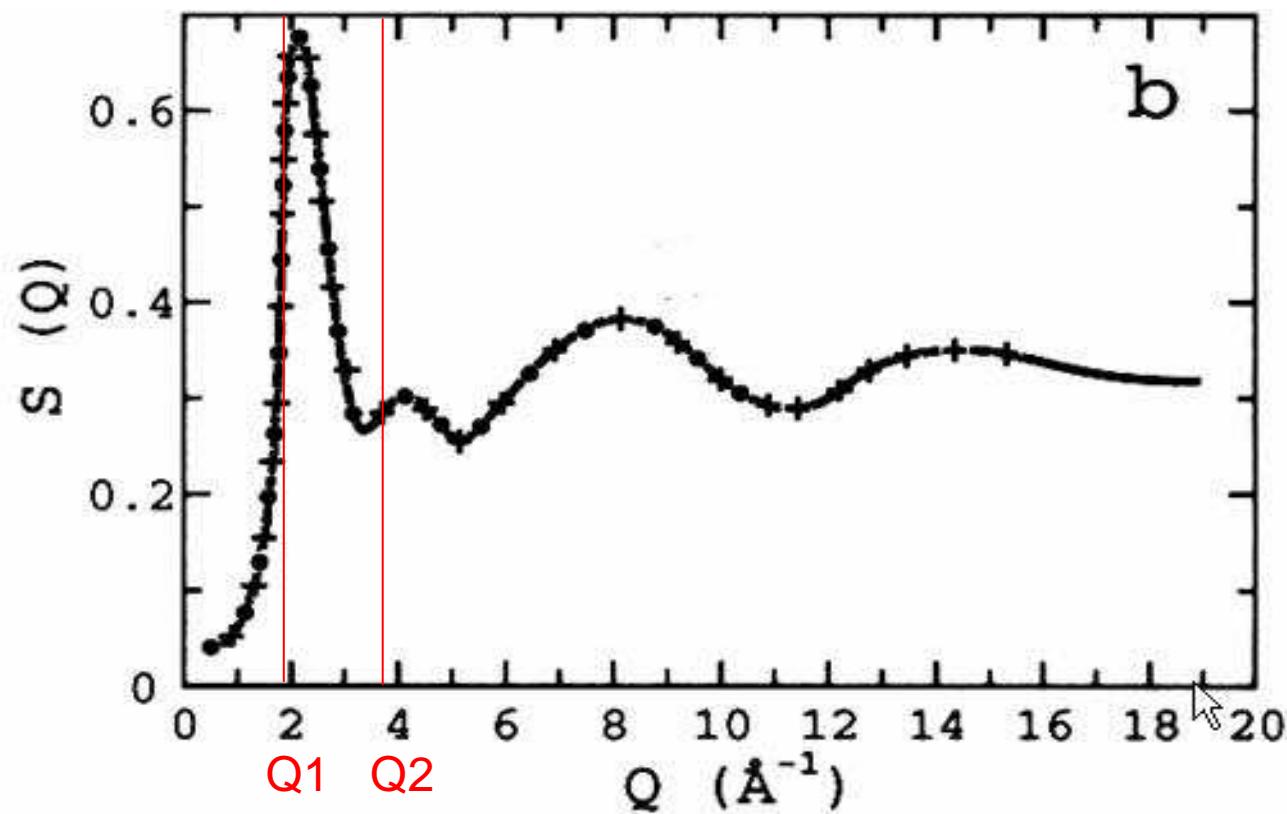
$S(Q)$ partials

A.K. Soper in *Hydrogen Bond Networks*
ed. M.-C. Bellissent-Funel and J. Dore
(Kluwer, Dordrecht, 1994), p. 97

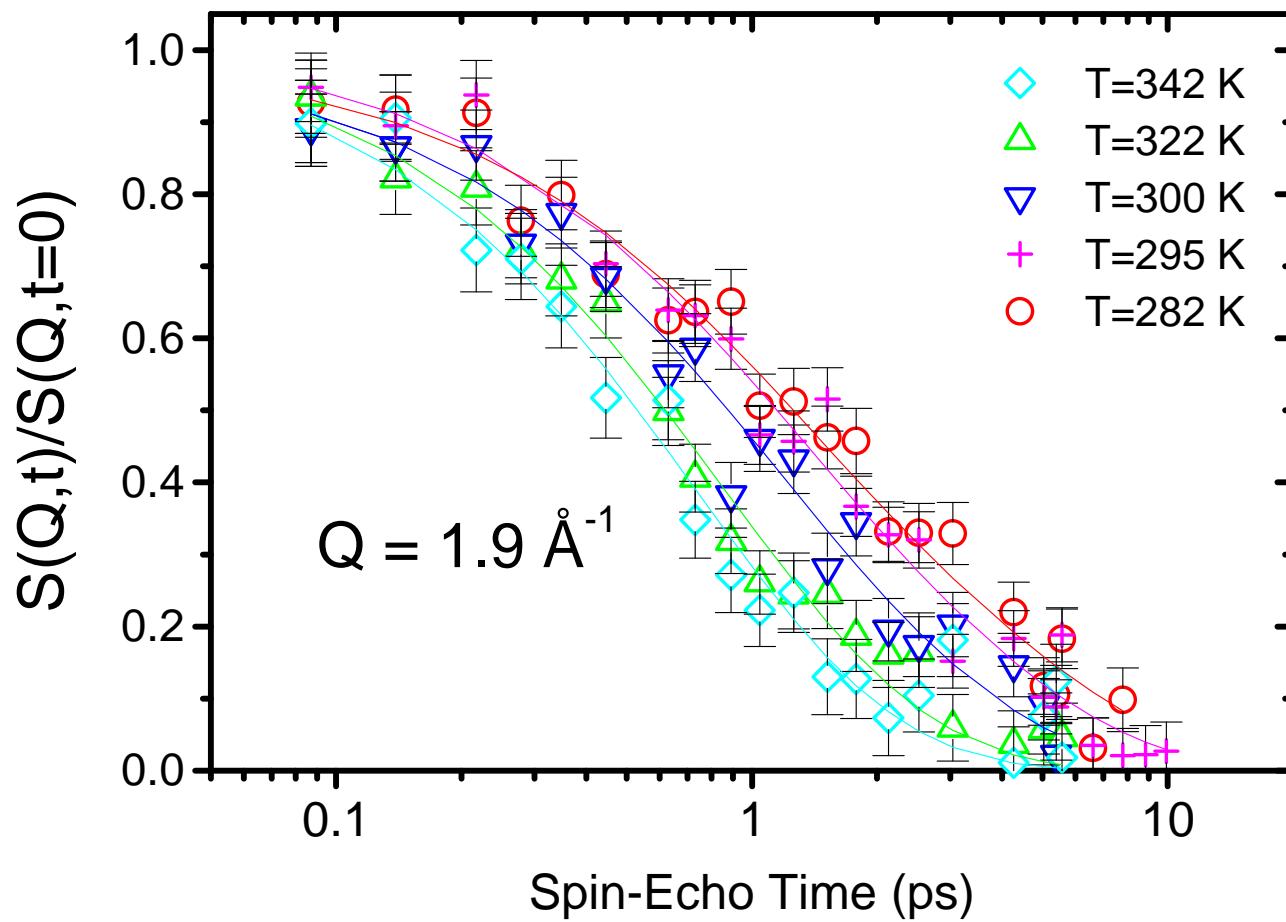
$$Q_1 = 1.9 \text{ \AA}^{-1}$$
$$Q_2 = 3.7 \text{ \AA}^{-1}$$



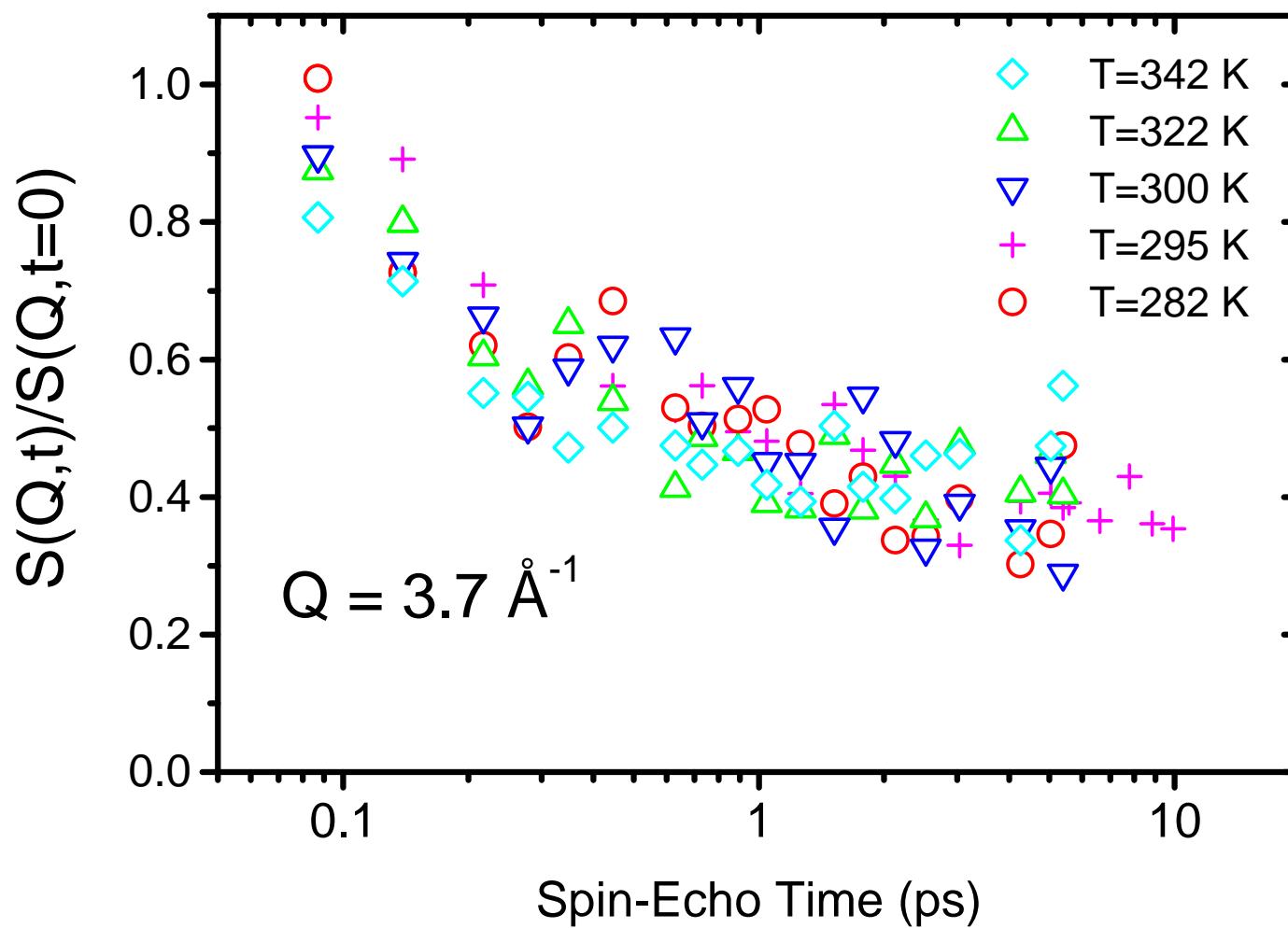
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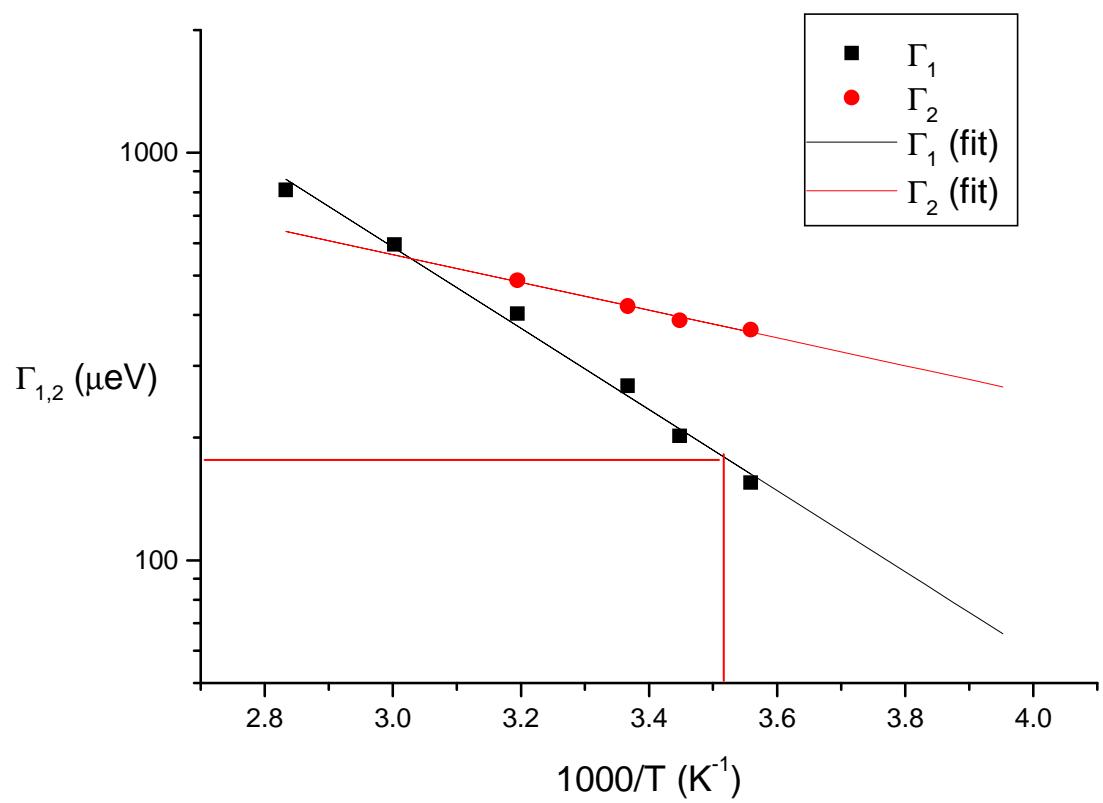
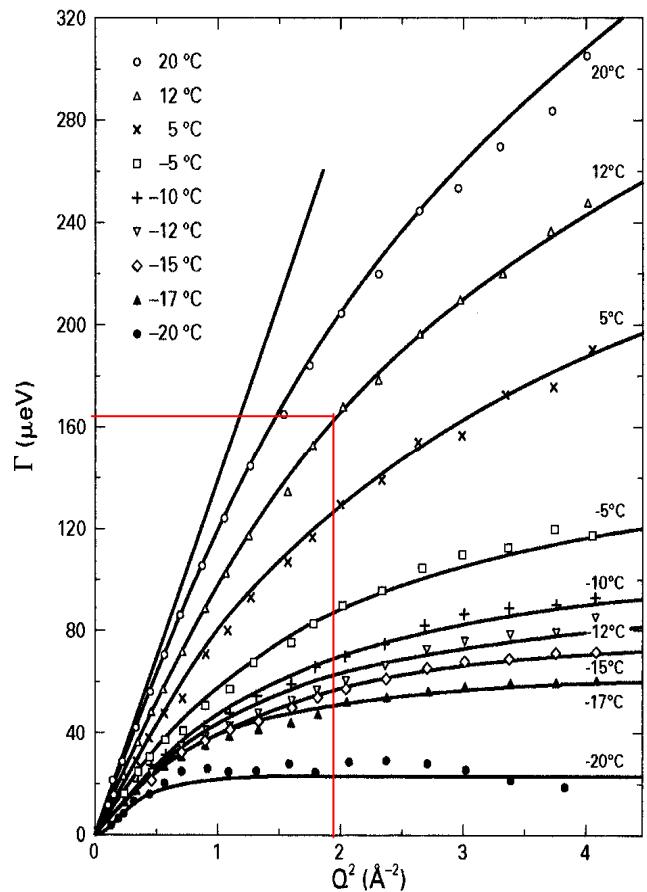


Intermediate scattering function

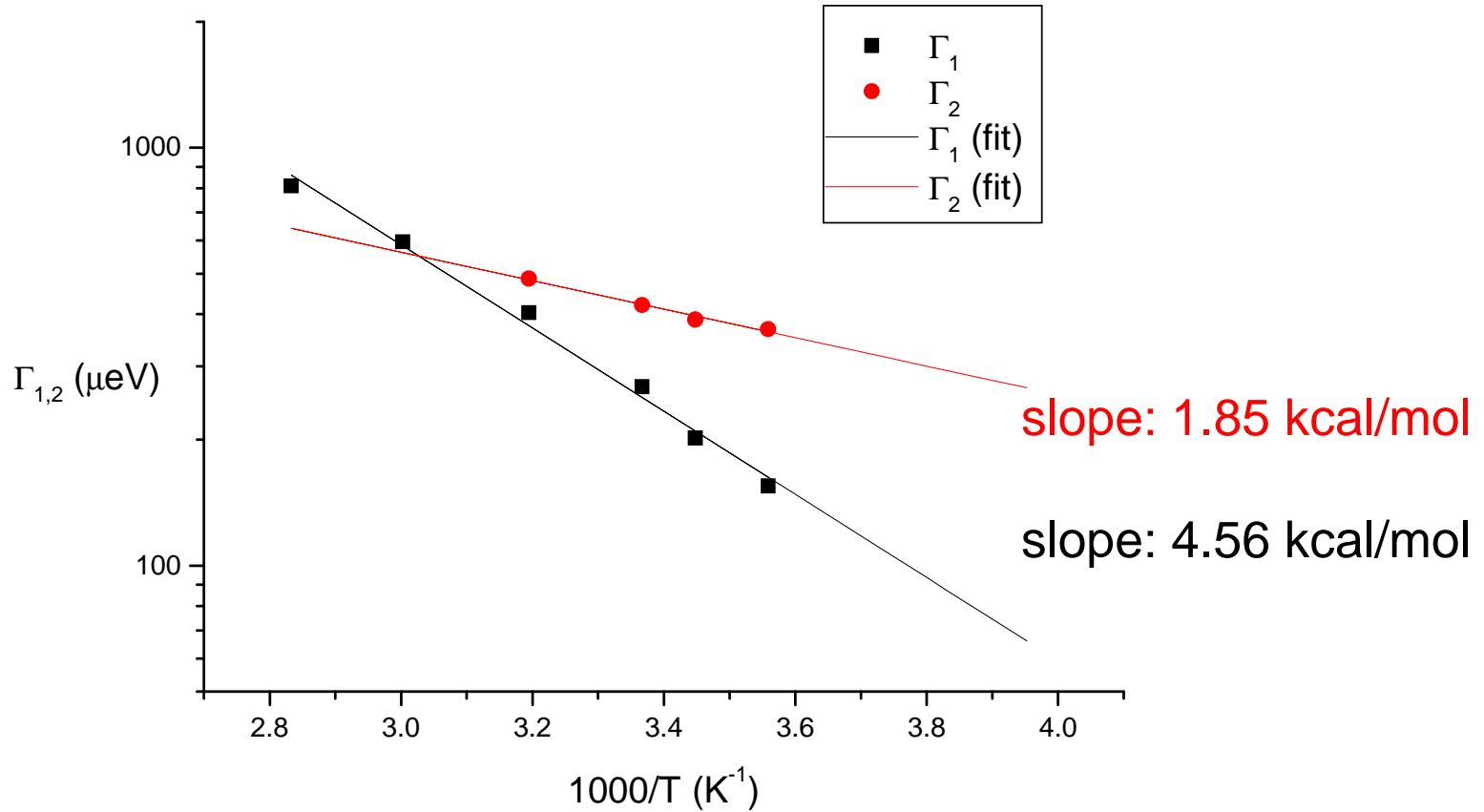


Intermediate scattering function





Line widths



$$\frac{4.56}{1.85} = 2.46 = n_{HB}$$

$$p_{HB} = n_{HB}/4 = 0.615$$

$$T = 296K = 23^\circ C$$

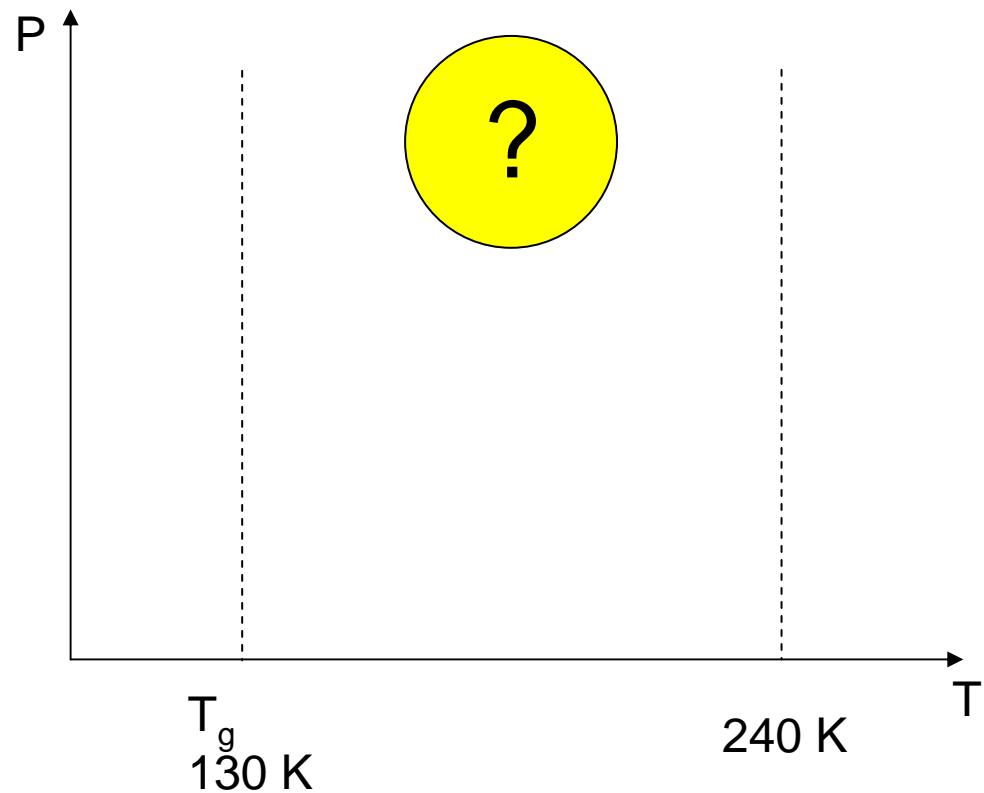
Experimental features

There is no discontinuity or divergence of any property of liquid water

Several independent experiments show that:

Hydrogen bond dynamics has an Arrhenius temperature dependence

It is possible to evaluate a "molecular time“ from the HB dynamics and from the number of intermolecular bonds. Its temperature dependence is strongly non-Arrhenius



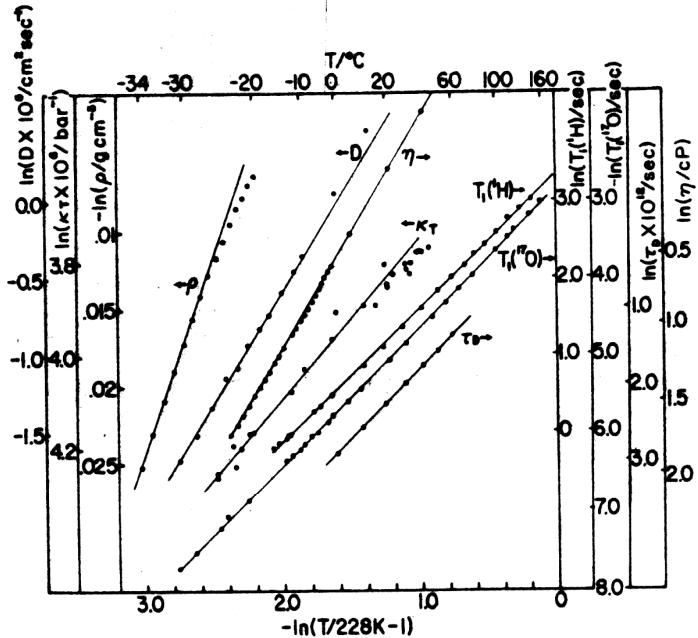


FIG. 4. $\ln W$ vs $\ln(T/T_s - 1)$ for the physical properties X of water at atmospheric pressure. The plots illustrate the applicability of Eq. (3) in the lowest temperature range for which data are available for each property. The temperature $T_s = 228$ K is used in each case although fractionally improved least squares best fits can be obtained with T_s up to 7°C removed from 228 K in individual cases. Data: ρ (density), (11, 12) D (diffusion coefficient), $^{15}\eta$ (viscosity), $^{16}\kappa_T$ (isothermal compressibility), $T_1(^1\text{H})$ (proton spin lattice relaxation time), $^{17}T_1(^{17}\text{O})$ (oxygen spin lattice relaxation time) $^{18}\tau_D$ (dielectric relaxation time). 19

R. Speedy and C.A. Angell
J. Chem. Phys. **65**, 851 (1976)

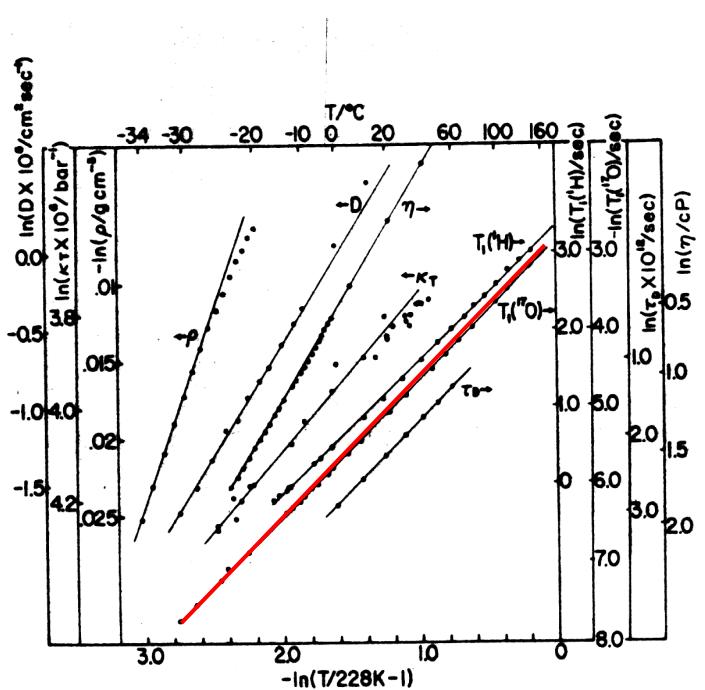
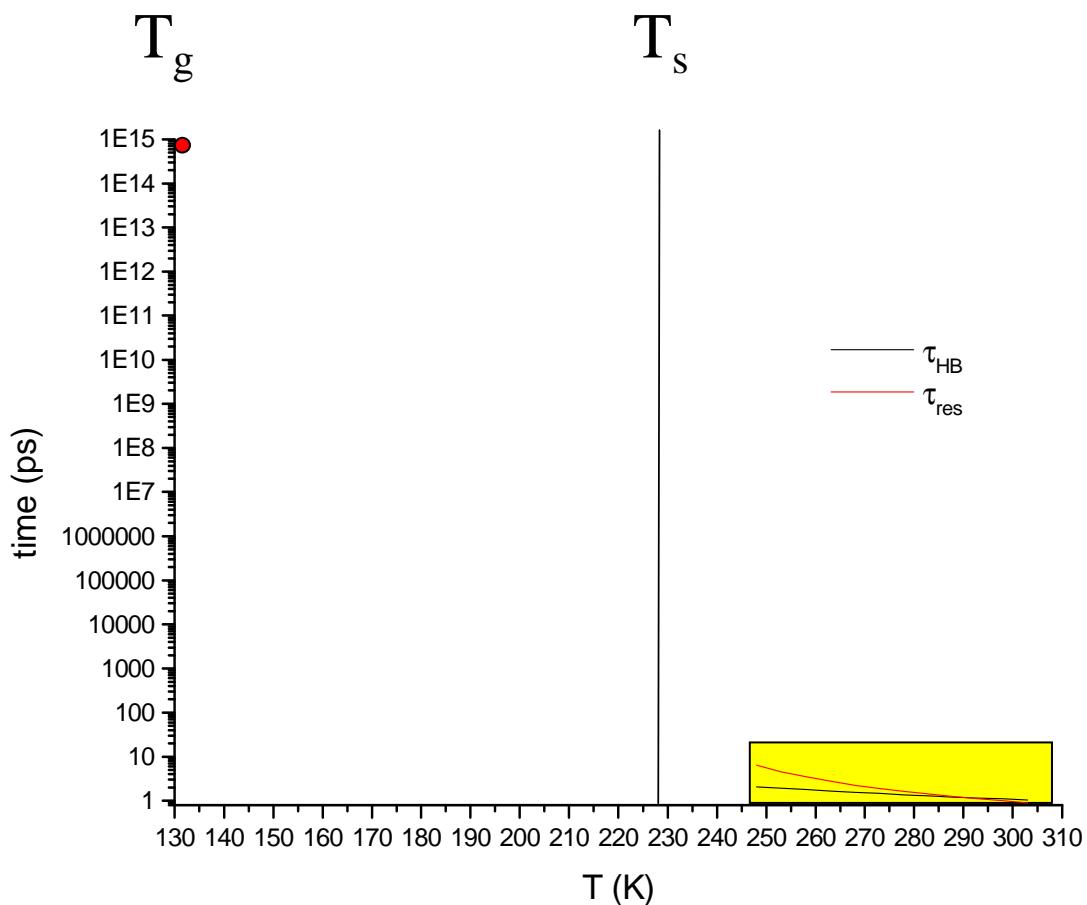
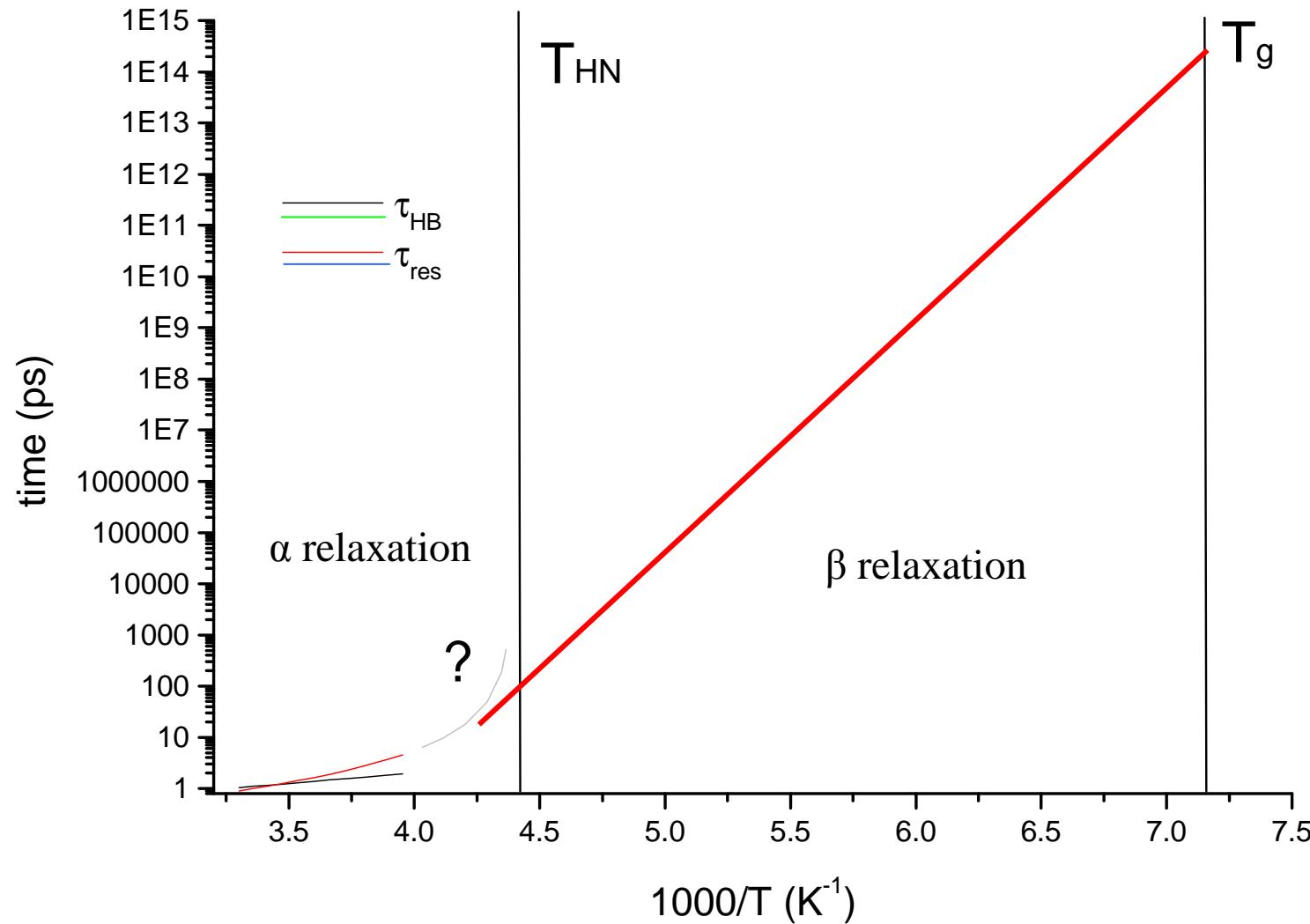


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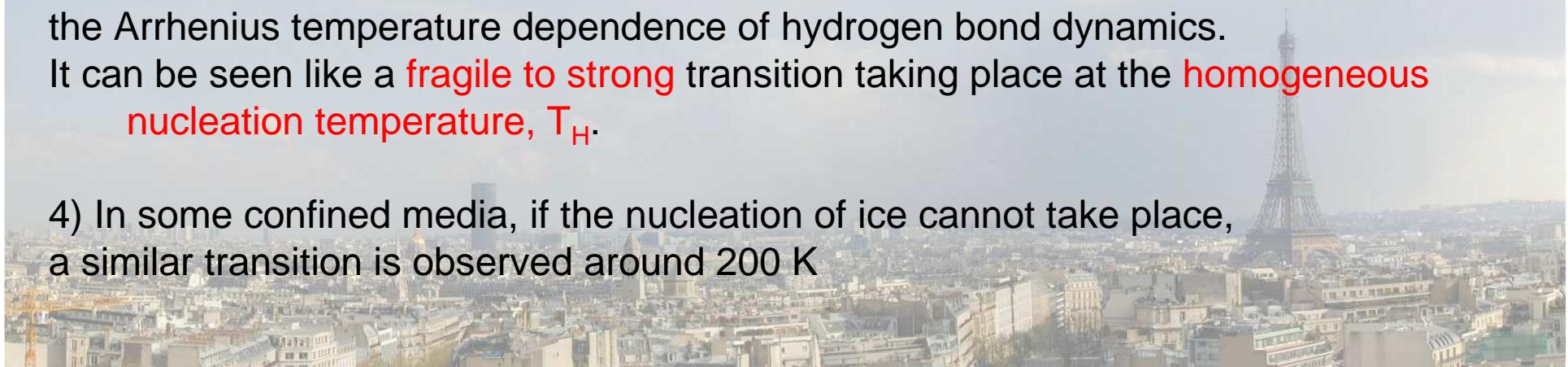
R. Speedy and C.A. Angell
J. Chem. Phys. **65**, 851 (1976)

A transition to Arrhenius temperature dependence?

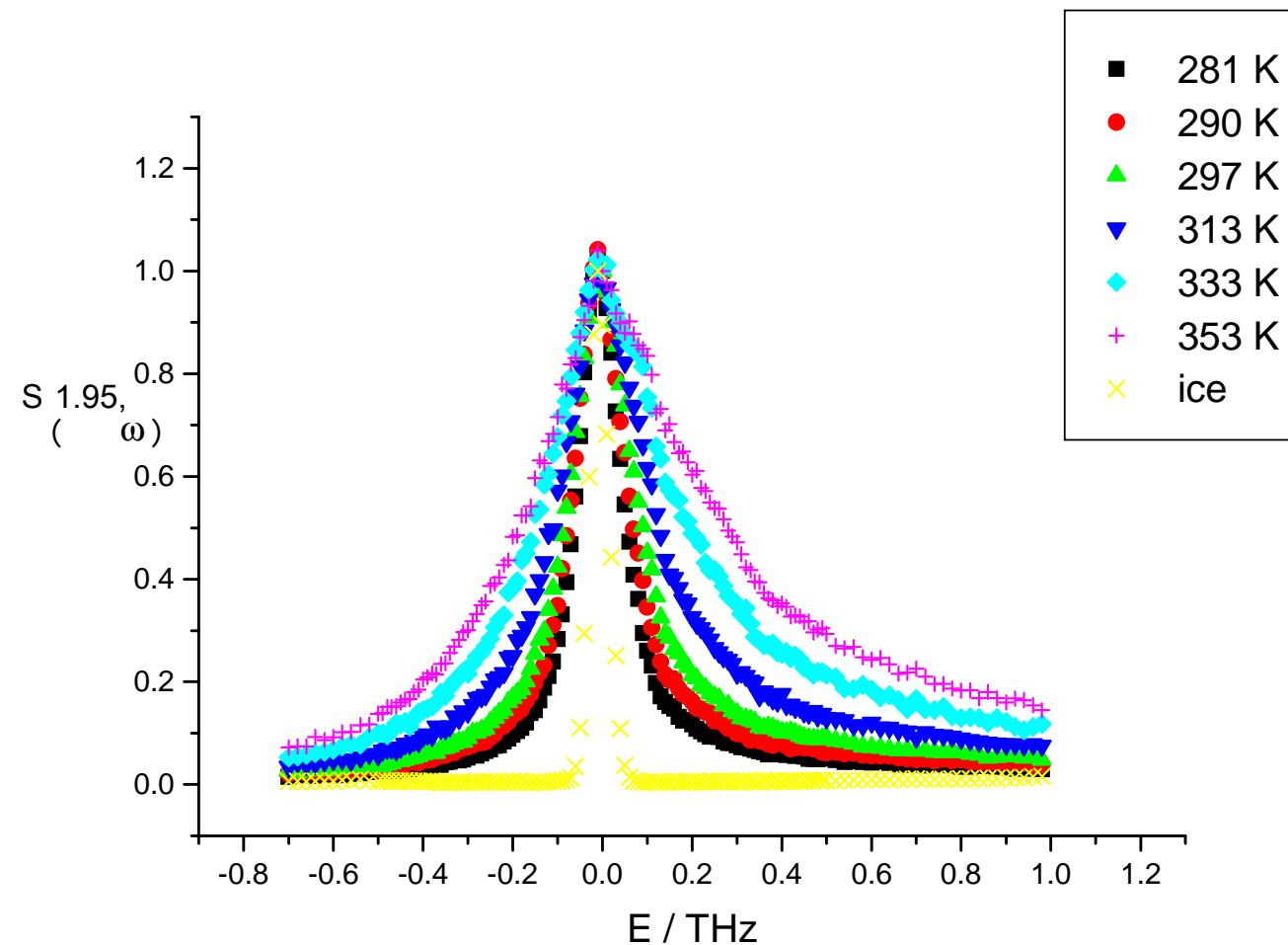


Conclusions

- 1) There are two main characteristic times :
 - a) A molecular residence time, τ_{res} , with a strong non-Arrhenius temperature dependence.
It shows an apparent divergence at the temperature of the homogeneous nucleation of hexagonal ice (-45 °C)
 - b) The hydrogen bond characteristic time, τ_{HB} , with an Arrhenius temperature dependence
- 2) The anomalous temperature dependence of τ_{res} can be predicted from τ_{HB} and the tetrahedral network of hydrogen bonds
- 3) During the quenching of bulk liquid water, if ice nucleation is avoided, there is a transition from a non-Arrhenius molecular behaviour to the Arrhenius temperature dependence of hydrogen bond dynamics.
It can be seen like a **fragile to strong** transition taking place at the **homogeneous nucleation temperature**, T_{H} .
- 4) In some confined media, if the nucleation of ice cannot take place, a similar transition is observed around 200 K

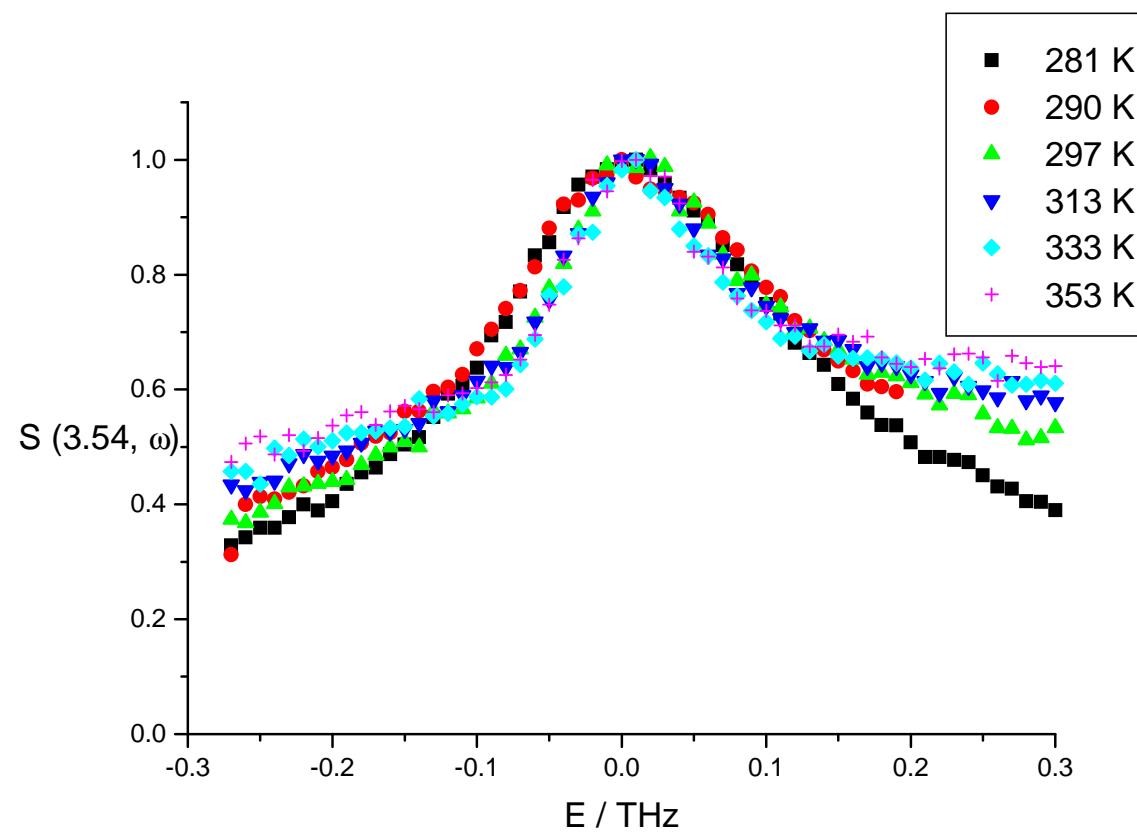


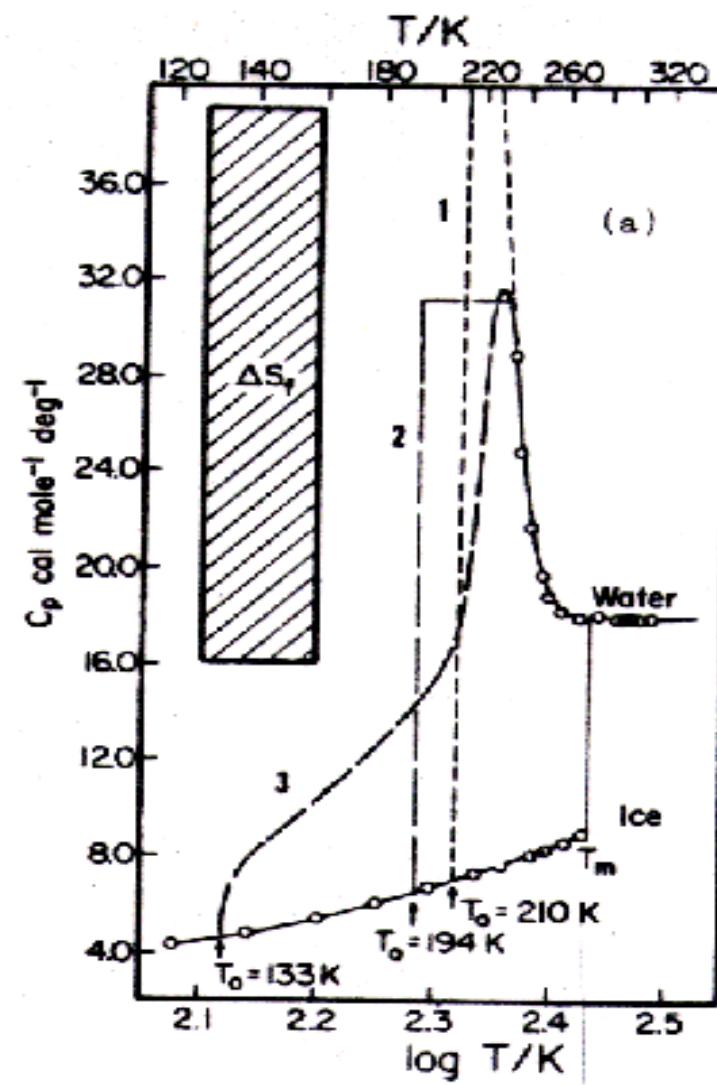
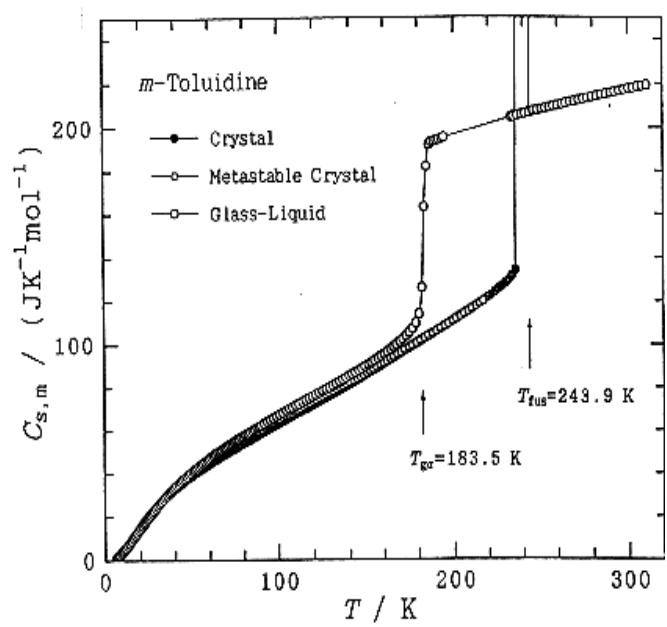
$S(Q, \omega)$ (D_2O)
 $Q = 1.95 \text{ \AA}^{-1}$



$S(Q, \omega)$ (D_2O)
 $Q = 3.54 \text{ \AA}^{-1}$

DD







$$S(Q) = (13.0 S_{\text{HH}}^s + 0.34 S_{\text{OO}}^s) \\ + (0.56 S_{\text{HH}}^d - 0.87 S_{\text{OH}}^d + 0.34 S_{\text{OO}}^d)$$



$$S(Q) = (1.22 S_{\text{DD}}^s + 0.34 S_{\text{OO}}^s) \\ + (1.78 S_{\text{DD}}^d + 1.55 S_{\text{OD}}^d + 0.34 S_{\text{OO}}^d)$$